TAMA Modecleaner alignment error signals Version 1.3

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1 Introduction

The effects of mirror misorientations^{[1](#page-0-0)} in interferometers with three or more mirrors cannot easily be determined analytically in the general case. Therefore a Mathematica program for numerical 3-D ray tracing was written. It uses geometrical optics to find the axis of the eigenmode for a given combination of reflecting and/or refracting plane and spherical surfaces. A similar program is described and printed in Appendix E3 of MPQ243. The term 'beam' in this text refers to the geometrical axis of a beam, without taking into account the transverse shape or optical phase.

2 The modecleaner cavities

Figure [1](#page-1-0) shows a schematic view of the TAMA modecleaner cavity seen from above, together with the coordinate system adopted in this section.

The cavity consists of two flat mirrors $(M_a \text{ and } M_b)$ that are separated by a relatively short distance (20 cm), and a curved mirror M_c with radius of curvature $R = 15$ m, which is at a distance $L = 9.738 \text{ m}$ from the (center between the) flat mirrors. The beam enters through M_a and travels clockwise to M_b , M_c , M_a , etc.

M^e is the location of the photodetector behind the end mirror (in the following called 'end detector' for short), which is at a distance of 2.08 m behind mirror M_c .

There are four beams of interest leaving the cavity. The main output beam is ' Out_b ' which goes to the interferometer and the stabilization of which is the main purpose of the modecleaner. There are two beams coming from M_a : the directly reflected input beam and the beam 'Outa' which is a fraction of the cavity eigenmode. The cavity is well aligned to the incoming beam (which we consider fixed), if these two beams are perfectly superimposed. By taking two quadrant diodes with two different lens systems and appropriately demodulating their outputs, we can obtain four independent error signals, similar to the case of a simple two-mirror Fabry-Perot cavity. The longitudinal locking signal is also obtained from these two interfering beams (using the Pound-Drever-Hall scheme).

 1 ¹To avoid confusion, we call a mirror or other component **misoriented** in this section, if its angular orientation differs from its reference orientation. The resulting movement of beams (e.g. cavity eigenmodes) will be called misalignment. An interferometer is called well-aligned if there are neither misorientations nor misalignments, i.e. components as well as beams are in their reference positions.

Figure 1: Schematic diagram of the TAMA modecleaner cavity seen from above.

In the three-mirror cavity, there are two additional degrees of freedom. They are linear combinations of movements of all three mirrors, as computed below. We will call them "neutral modes" because they have no influence on either the interference of the incoming beam with the cavity eigenmode nor on the outgoing beam. They can be used to control the spot position on the far mirror M_c . This is described in more detail below.

In the ray-tracing program, we first compute the well-aligned case (i.e., all mirrors are hit in their center) as reference. We call P_a , P_b and P_c the points where the axis of the eigenmode intersects the mirrors M_a , M_b and M_c , respectively. After misorienting one particular mirror by the small angle ε , we recompute the eigenmode axis, compare it with the well-aligned case and divide the difference by ε . The main results are:

- The shifts of the spots P_a , P_b and P_c .
- The angles γ_a , γ_b and γ_c between the beams 'Out_a', 'Out_b', 'Out_c' and their respective references. For the vertical misalignments which are considered separately, we call these angles δ_a , δ_b and δ_c , respectively. We also compute the angle γ_d (δ_d) between the directly reflected incoming beam and its reference for the case that M^a is misoriented.
- For 'Out_a' in the case of misorienting M_a , we also compute the angle $\gamma'_a = \gamma_a \gamma_d$, which is the angle between the beam leaving the cavity and the directly reflected beam, because this is the angle between the interfering wavefronts that is detected by the quadrant diode.
- As described above we finally compute the angle θ which describes the 'character' of the misalignment at the waist. It is given by $\theta^w = \arctan(\gamma'_a z_R/\Delta z_{waist})$.

The waist of the cavity eigenmode is located halfway between the mirrors M_a and

Mb. Its Rayleigh range is given by

$$
z_{\rm R} = \sqrt{\frac{L_{\rm RT}}{2} \left(R - \frac{L_{\rm RT}}{2} \right)} = 7.126 \,\mathrm{m},\tag{1}
$$

where $L_{RT}/2$ is one half of the round-trip distance:

$$
L_{RT}/2 = \sqrt{L^2 + d^2/4} + d/2 = 9.8385 \,\text{m}.\tag{2}
$$

3 Horizontal misalignments:

By 'horizontal' misalignments we mean that a mirror is rotated around the vertical axis, i.e. beam spots move horizontally (in the plane of the modecleaner cavity). Angles are counted as positive if a mirror is rotated clockwise, if seen from above (as in Figure [1\)](#page-1-0).

We introduce the common- and differential mode motion of mirrors M_a and M_b by defining angles α_+ and α_- . Furthermore we introduce the "neutral" mode α_n as follows:

$$
\begin{array}{c|c|c}\n & M_{\rm a} & M_{\rm b} & M_{\rm c} \\
\hline\n\alpha_{+} & \alpha_{\rm a} = \alpha_{+} & \alpha_{\rm b} = \alpha_{+} & \alpha_{\rm c} = 0 \\
\alpha_{-} & \alpha_{\rm a} = \alpha_{-} & \alpha_{\rm b} = -\alpha_{-} & \alpha_{\rm c} = 0 \\
\alpha_{\rm n} & \alpha_{\rm a} = \alpha_{\rm n} & \alpha_{\rm b} = \alpha_{\rm n} & \alpha_{\rm c} = 0.701\alpha_{\rm n}\n\end{array} \tag{3}
$$

$$
\alpha_{+} = \frac{\alpha_{a} + \alpha_{b}}{2},\tag{4}
$$

$$
\alpha_{-} = \frac{\alpha_{\rm a} - \alpha_{\rm b}}{2},\tag{5}
$$

The results of the raytracing program are given in Table [1.](#page-2-0)

	Cause	P_a		$P_{\rm a}$		waist	Out_a		Out'_2		$\theta^{\rm w}$	
			Δx	Δz		Δz		$\gamma_{\rm a}$		$\gamma'_{\rm a} = \gamma_{\rm a} - \gamma_{\rm d}$		
	$\alpha_{\rm a}$	-9.740 m \cdot α _a	$9.840 \,\mathrm{m} \cdot \alpha_{a}$			$9.739 \,\mathrm{m} \cdot \alpha_{\mathrm{a}}$		$0.981\alpha_{\rm a}$		$-1.019 \alpha_{\rm a}$	-36.72°	
	$\alpha_{\rm b}$	$9.538 \text{ m} \cdot \alpha_{\text{h}}$	-9.636 m \cdot α _b			-9.739 m \cdot α_{h}		$-1.019\alpha_{h}$		$-1.019\alpha_{\rm h}$	36.72°	
α_c		$0.288 \,\mathrm{m} \cdot \alpha_c$ -0.291 m $\cdot \alpha_c$				$0.000 \,\mathrm{m} \cdot \alpha_c$	$2.906 \alpha_c$		$2.906\,\alpha_c$		-90.00°	
α_+		-0.202 m $\cdot \alpha_+$	$0.204 \,\mathrm{m}\cdot\alpha_{+}$			$0.000 \text{ m} \cdot \alpha_+$	$-0.039 \alpha_+$		$-2.039 \alpha_+$		90.00°	
	α —	-19.278 m $\cdot \alpha$ –	$19.477 \,\mathrm{m} \cdot \alpha$		$19.477 \,\mathrm{m} \cdot \alpha$			2.000α –		0.000α –	0.00°	
	$\alpha_{\rm n}$		$0.000 \,\mathrm{m} \cdot \alpha_{\mathrm{n}}$ $0.000 \,\mathrm{m} \cdot \alpha_{\mathrm{n}}$		$0.000 \,\mathrm{m} \cdot \alpha_{\mathrm{n}}$			$2.000\,\alpha_n$		$0.000 \alpha_n$		
Cause		P _b	P _b		Out _b		P_c		Out _c	d. refl.		End
		Δx	Δz				Δx					Δx
					$\gamma_{\rm b}$				$\gamma_{\rm c}$	γ_{d}		
α_{a}		$9.538 \text{ m} \cdot \alpha_{\text{a}}$	$9.637 \text{ m} \cdot \alpha_{\text{a}}$	$1.019 \alpha_{\rm a}$		-0.291 m \cdot α _a		$-1.019\alpha_{\rm a}$		$2 \cdot \alpha_{\rm a}$	$-2.411 \text{ m} \cdot \alpha_{\rm a}$	
$\alpha_{\rm b}$		$-9.740 \,\mathrm{m} \cdot \alpha_{\mathrm{b}}$	$-9.841 \,\mathrm{m} \cdot \alpha_{\mathrm{b}}$	$1.019\,\alpha_{\rm b}$		$-0.291 \text{ m} \cdot \alpha_{\text{h}}$		$0.981\alpha_{\rm h}$		$0 \cdot \alpha_{\rm b}$	$1.749 \,\mathrm{m} \cdot \alpha_{\mathrm{b}}$	
α_c		$0.288 \,\mathrm{m} \cdot \alpha_c$	$0.291 \,\mathrm{m} \cdot \alpha_c$	$-2.906 \alpha_c$		$28.596 \,\mathrm{m} \cdot \alpha_c$		$2.906 \alpha_c$		$0 \cdot \alpha_c$	$34.642 \,\mathrm{m} \cdot \alpha_c$	
α_+		$-0.202 \,\mathrm{m} \cdot \alpha_+$	-0.204 m $\cdot \alpha_+$	$2.039\,\alpha_{+}$		-0.581 m $\cdot \alpha_+$		$-0.039 \alpha_{+}$		$2 \cdot \alpha_+$	-0.662 m $\cdot \alpha_+$	
α $-$		$19.278 \,\mathrm{m} \cdot \alpha$	19.477 m $\cdot \alpha$	0.000α		$0.000 \,\mathrm{m} \cdot \alpha$		-2.000α		$2 \cdot \alpha$	$-4.160 \,\mathrm{m} \cdot \alpha$	
$\alpha_{\rm n}$	$0.000 \,\mathrm{m} \cdot \alpha_{\mathrm{n}}$		$0.000 \,\mathrm{m} \cdot \alpha_{\mathrm{n}}$	$0.000\alpha_n$		$19.478\,\mathrm{m}\cdot\alpha_{\mathrm{n}}$		$2.000\alpha_n$		$2 \cdot \alpha_{\rm n}$	$23.639\,\mathrm{m}\cdot\alpha_{\mathrm{n}}$	

Table 1: Results of the ray-tracing program for horizontal misalignments of the TAMA modecleaner.

4 Vertical misalignments:

In the modecleaners, the horizontal and vertical axes are not similar. The results of the raytracing program for vertical misalignments are given in Table [2.](#page-4-0) Note, for example, that a small vertical tilt β_a of the input mirror M_a (which is hit under approximately 45° from the incoming beam) causes a deflection of the reflected beam by only $\delta_d = 1.43\beta_a$ as compared to $\gamma_d = 2\alpha_a$ in the horizontal case. Another example is the tilt of M_c which, if horizontal, causes a pure angular misalignment at the waist. A vertical tilt of M_c , on the other hand, shifts the cavity eigenmode downwards parallely, without changing any angles. Angles are now counted as positive when the mirror normal moves downward from the reference direction.

As before we introduce the common- and differential mode motion of mirrors M_a and M_b by defining angles β_+ and β_- . Furthermore we introduce the "neutral" mode β_n as follows:

$$
\begin{array}{|c|c|c|c|c|}\n\hline\n & M_{\rm a} & M_{\rm b} & M_{\rm c} \\
\hline\n\beta_{+} & \beta_{\rm a} = \beta_{+} & \beta_{\rm b} = \beta_{+} & \beta_{\rm c} = 0 \\
\beta_{-} & \beta_{\rm a} = \beta_{-} & \beta_{\rm c} = 0 & \beta_{\rm n} \\
\beta_{\rm n} & \beta_{\rm a} = \beta_{\rm n} & \beta_{\rm b} = \beta_{\rm n} & \beta_{\rm c} = -0.4986 \beta_{\rm n}\n\end{array} \tag{6}
$$

$$
\beta_{+} = \frac{\beta_{\rm a} + \beta_{\rm b}}{2},\tag{7}
$$

$$
\beta_- = \frac{\beta_a - \beta_b}{2},\tag{8}
$$

The results of the raytracing program are given in Table [2.](#page-4-0)

5 Degrees of freedom

The most important alignment task is to superimpose the axis of the cavity eigenmode with the axis of the incoming beam. This requires the control of four degrees of freedom. For this purpose, in the differential wavefront sensing method, we place two quadrant detectors with different lens systems in the beam reflected from Ma. The interference between the directly reflected incoming beam, which is phase modulated at an RF frequency, and the beam 'Outa' leaking out of the cavity contains enough information to lock the cavity longitudinally and to obtain alignment error signals for those four degrees of freedom that determine the superposition of the incoming beam and the cavity eigenmode.

In particular, we now assume all mirrors to be slightly misoriented and compute the combined signals which are obtained by demodulating the outputs of two quadrant detectors, one (called X_I) with $\Phi = 0^\circ$ and the other one (called X_O) with $\Phi = 90^\circ$ of extra phase shift. We scale parallel shifts Δy or Δz with the appropriate factor z_R and obtain for horizontal misalignments:

$$
X_I = -1.019 \alpha_a - 1.019 \alpha_b + 2.906 \alpha_c = -1.019 \alpha_+ + 2.906 \alpha_c, \tag{9}
$$

$$
X_Q = 1.366 \,\alpha_a - 1.366 \,\alpha_b = 1.366 \,\alpha_-, \tag{10}
$$

Cause	$P_{\rm b}$	Out _b	P_c	Out _c	d. refl.	End
	Δy	$\delta_{\rm b}$	Δy	δ_c	$\delta_{\rm d}$	Δy
$\beta_{\rm a}$	$-3.810 \,\mathrm{m} \cdot \beta_{\rm a}$	$0.704\,\beta_{\rm a}$	$-10.661 \,\mathrm{m} \cdot \beta_{\rm a}$	$0.704\,\beta$ _a	$1.42 \cdot \beta_{\rm a}$	$-12.125 \,\mathrm{m} \cdot \beta_{\rm a}$
$\beta_{\rm b}$	$-3.670 \,\mathrm{m} \cdot \beta_{\mathrm{b}}$	$-0.704\beta_{\rm h}$	$-10.661 \,\mathrm{m} \cdot \beta_{\mathrm{b}}$	$0.718 \beta_{\rm h}$	$0.00 \cdot \beta_{\rm b}$	$-12.155 \,\mathrm{m} \cdot \beta_{\mathrm{b}}$
$\beta_{\rm c}$	$-15.000 \,\mathrm{m} \cdot \beta_c$	$0.000\,\beta_{\rm c}$	$-15.000 \,\mathrm{m} \cdot \beta_c$	$0.000\,\beta_{\rm c}$	$0 \cdot \beta_c$	$-15.000 \,\mathrm{m} \cdot \beta_c$
β_+	$-7.480\,\mathrm{m}\cdot\beta_+$	$0.000\,\beta_{+}$	$-21.323 \,\mathrm{m} \cdot \beta_+$	$1.421 \,\beta_{+}$	$-1.421 \cdot \beta_{+}$	$-24.280 \,\mathrm{m} \cdot \beta_+$
β_{-}	-0.141 m \cdot β	$1.407\,\beta$	$0.000 \,\text{m} \cdot \beta_-$	$-0.014\,\beta$	$-1.421 \cdot \beta_{-}$	$0.030\,\text{m}\cdot\beta$
$\beta_{\rm n}$	$0.000 \,\mathrm{m} \cdot \beta_{\rm n}$	$0.000 \beta_{\rm n}$	$-13.843 \,\mathrm{m} \cdot \beta_{\rm n}$	$1.421\,\beta_{\rm n}$	$-1.420 \cdot \beta_n$	$-16.800 \,\mathrm{m} \cdot \beta_n$

Table 2: Results of the ray-tracing program for vertical misalignments of the TAMA modecleaner.

and for vertical misalignments:

$$
Y_I = 0.7035 \,\beta_a - 0.7035 \,\beta_b = 0.7035 \,\beta_-, \tag{11}
$$

$$
Y_Q = -0.524 \,\beta_a - 0.524 \,\beta_b - 2.105 \,\beta_c = -0.524 \,\beta_+ - 2.105 \,\beta_c. \tag{12}
$$

Using simple linear algebra we also find the "neutral" modes from these results. They are given in tables [\(3\)](#page-2-1) and [\(6\)](#page-3-0) above.

By inspecting tables [1](#page-2-0) and [2](#page-4-0) one finds that the main effect of these "neutral modes" is to shift the spot position on the far end mirror Mc. Hence their control is not absolutely essential for the basic function of the modecleaner. In the present situation, there is neither feedback to M_c nor is the spot position on M_c monitored. One problem with this approach might be that according to equations (9) to (12) , a small motion of M^c translates into error signals at the reflected light port which is 3. . . 4 times larger than the error signal for a motion of the front mirrors of comparable amplitued. This huge error signal is then fed back the the *front* mirrors which are forced to compensate Mc's misorientation although they are not very efficient actuators for that purpose. Furthermore, such feedback shifts the spot position on M_c which due to the curved mirror surface again causes the whole circle to start. It is possible that some of the present dynamic range problems may be relieved by sensing and feeding back all degrees of freedom.

Because the main effect of the "neutral modes" is to shift the spot position on the far end mirror M_c , the most straightforward approach is to sense this spot position (with a DC quadrant detector looking at the transmitted light). In tables [1](#page-2-0) and [2](#page-4-0) above, this shift is computed as Δx or Δy of point Pc, respectively.

However, the detector is located at a distance $e = 2.08 \text{ m}$ behind M_c , and the spot

position detected is hence a linear combination^{[2](#page-5-0)} of Δx or Δy of point Pc, respectively with the angle δ_{c} .

This shift of the beam spot on the photodetector is also calculated by the program and printed as "End Δy " in Tables [1](#page-2-0) and [2.](#page-4-0)

If we also scale it by the factor z_R for consistency (although the way of detection is different and hence an arbitrary scale factor appears anyway), and call the error signals thus obtained X_E and Y_E , we get

$$
X_E = -0.3383 \alpha_a + 0.2455 \alpha_b + 4.8613 \alpha_c \tag{13}
$$

$$
= -0.0464 \alpha_{+} - 0.2919 \alpha_{-} + 4.8613 \alpha_{c}, \tag{14}
$$

$$
Y_E = -1.701 \,\beta_a - 1.705 \,\beta_b - 2.105 \,\beta_c \tag{15}
$$

$$
= -1.703 \beta_{+} + 0.002 \beta_{-} - 2.105 \beta_{c}.
$$
\n(16)

We see that for the horizontal direction, the additional error signal is dominated by the misorientation of M_c and could directly be fed back to that mirror, while for the vertical direction the signals are more mixed. However, they are still sufficiently linearly independent such that by inversion of the well-conditioned matrix error signals for all three mirrors can be found. This matrix inversion yields

$$
\alpha_{\rm a} = X_E -1.6726 X_I +1.4273 X_Q \n\alpha_{\rm b} = X_E -1.6726 X_I -X_Q \n\alpha_{\rm c} = 0.7015 X_E -0.0319 X_I +0.1498 X_Q
$$
\n(17)

$$
\begin{array}{rcl}\n\beta_{a} & = & -Y_{E} & +1.6785 \, Y_{I} & +Y_{Q} \\
\beta_{b} & = & -Y_{E} & -1.6785 \, Y_{I} & +Y_{Q} \\
\beta_{c} & = & 0.4987 \, Y_{E} & -0.0015 \, Y_{I} & -1.6186 \, Y_{Q}\n\end{array} \tag{18}
$$

where scaling factors of $1/0.301462$ and $1/0.42418$, respectively, have been applied to make the most common coefficients unity. Note that the present alignment system uses only a 2×2 matrix without X_E , Y_E , α_c and β_c .

6 Actuator transfer functions

The transfer functions from all actuators to all sensors have been measured by K. Arai. The results were fitted with LISO and are shown on the following pages.

²Many thanks to K. Arai for noting this point.

6.1 Input Mirror Yaw coil ("m1xcoil")

 α

pole0:f = 1.372388 pole0:q = 0.755353 pole1:f = 3.777486

- pole1:q = 3.455169 zero0:f = 2.508392
- zero0:q = 0.840509

6.2 Output Mirror Yaw coil ("m2xcoil")

20

pole0:q = 0.804331 pole1:f = 3.661688 pole1:q = 3.842627 zero0:f = 2.316205 $zero0:q = 0.669030$

6.3 Input Mirror Pitch PZT ("m1ypzt")

pole1:f = 20.08673

pole1:q = 16.34634

6.4 Output Mirror Pitch PZT ("m2ypzt")

10

pole0:q = 11.15256 pole1:f = 20.48184 pole1:q = 28.21300

0

pole0:f = 8.757994 pole0:q = 11.90231 pole1:f = 20.31437 pole1:q = 27.83178 zero0:f = 14.49367

zero0:q = 12.04289

6.6 Output Mirror Pitch coil ("m2ycoil")

10

 $pole0:q = 11.66273$ pole1:f = 20.58340 pole1:q = 26.94568 zero0:f = 14.29287 zero0:q = 10.79691

6.7 End Mirror Yaw Coil ("m3xcoil")

pole0:f = 1.8826161 pole0:q = 1.3989960

 Ω

pole0:f = 6.436784 pole0:q = 8.052757 pole1:f = 15.49933 pole1:q = 10.58777

6.9 End Mirror Pitch coil ("m3ycoil")

0

pole0:q = 11.48701 pole1:f = 15.65282 pole1:q = 16.45207 zero0:f = 13.19353 $zero0:q = 9.075651$

7 Calibration Factors

Both actuators and sensors have calibration factors which are not yet known. In order to find them (and test the above theroretical model), the LISO-fitted factors from the transfer functions were fitted to simple model

$$
factor = A_{\text{optical}} \times G_{\text{actuator}} \times G_{\text{sensor}}.\tag{19}
$$

where *all* measurements were fitted simultaneously to one set of calibration constants G_i .

For that purpose a simple C program ("prod.c") was written which uses the Nelder-Mead Simplex algorithm to minimize

$$
\chi^2(G_1, \dots, G_m) = \sum_{i=1}^{n_{\text{data}}} w_i \left| f_i - \tilde{f}_i(G_1, \dots, G_m) \right|^2 \tag{20}
$$

where f_i are the measured factors, \tilde{f}_i their approximations from Equation [19,](#page-15-0) $w_i =$ $1/\sigma_i^2$ are the weights of each data point derived from their standard deviation (taken from the LISO fit), and G_i the unknown calibration constants.

(21)

The lines marked *** at the end were excluded from the fit, because the theoretically expected factor was zero (and the measured one was indeed small).

The resulting calibration constants are:

The residuals were:

(22)

(23)